

Killing fields and conservation laws

Daniel Fox

June 13, 2011

Goals

- Illustrate a connection between

Killing Fields of primitive maps $dB_\lambda + [\psi_\lambda, B_\lambda] = 0$	Conservation Laws as elements of the characteristic cohomology $\varphi \in \bar{H}^1(M^{(\infty)})$
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Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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- This connection stems from the work of Terng and Wang (2004), Pinkall and Sterling (1989), and Burstall, Ferus, Pedit, and Pinkall (1993)—The point is to incorporate some of these techniques into the framework of the characteristic cohomology

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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- Use the connection above to completely determine the conservation laws as elements of the characteristic cohomology for

$$u_{z\bar{z}} = -f(u) \quad \text{where } u : \mathbb{C} \rightarrow \mathbb{R}.$$

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

- Reinterpret **finite type solutions** using conservation laws

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

- Reinterpret **finite type solutions** using conservation laws
- Raise questions about what information is stored in conservation laws, thought of as cohomology classes

Nonlinear Poisson equation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

In order to encode

$$\frac{\partial^2 u}{\partial z \partial \bar{z}} = -f(u),$$

as an EDS with independence condition, let

$M = J^1(\mathbb{C}, \mathbb{R}) = \mathbb{C} \times \mathbb{R} \times \mathbb{C}$ have coordinates (z, u, u_0) and define the differential forms

$$\zeta = dz$$

$$\eta_0 = du - u_0 \zeta - \bar{u}_0 \bar{\zeta}$$

$$\psi = -\frac{\sqrt{-1}}{2}(\zeta \wedge du_0 - \bar{\zeta} \wedge d\bar{u}_0 + 2f dz \wedge d\bar{z}).$$

Nonlinear Poisson equation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathcal{I} = \langle \eta_0, \psi \rangle.$$

Nonlinear Poisson equation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Independence condition: $\zeta \wedge \bar{\zeta} \neq 0$.

Nonlinear Poisson equation: infinite prolongation

To study higher-order conservation laws we turn to the infinite prolongation $M^{(\infty)}$

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$$\{z, u, u_0, u_1, \dots\}$$

Nonlinear Poisson equation: infinite prolongation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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where $u_i = \frac{\partial^{i+1}}{\partial z^{i+1}} u$ on solutions.

Nonlinear Poisson equation: infinite prolongation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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If $A : M^{(\infty)} \rightarrow \mathbb{C}$ then $dA \equiv e_{-1}(A)\zeta + \overline{e_{-1}}(A)\overline{\zeta}$ modulo $I^{(\infty)}$.

Nonlinear Poisson equation: \mathbb{S}^1 -symmetry and weighted-degree

$$(z, u) \rightarrow (\lambda^{-1}z, u)$$

is a symmetry of

$$u_{z\bar{z}} = -f(u)$$

for $\lambda \in \mathbb{C}$, $|\lambda| = 1$.

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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\mathbb{S}^1 -action on $M^{(\infty)}$:

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$F_\lambda : M^{(\infty)} \rightarrow M^{(\infty)}$$

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$F_\lambda : M^{(\infty)} \rightarrow M^{(\infty)}$$

$$F_\lambda(z, u, u_j) = (\lambda^{-1}z, u, \lambda^{j+1}u_j)$$

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$F_\lambda^* I^{(\infty)} = I^{(\infty)}$$

Conservation laws as elements of the characteristic cohomology

Let $(M^{(\infty)}, \mathcal{I}^{(\infty)})$ be the infinite prolongation of the EDS corresponding to $u_{z\bar{z}} = -f$.

Conservation laws as elements of the characteristic cohomology

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Group of Conservation Laws:

Conservation laws as elements of the characteristic cohomology

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Let $(M^{(\infty)}, \mathcal{I}^{(\infty)})$ be the infinite prolongation of the EDS corresponding to $u_{z\bar{z}} = -f$.

Group of Conservation Laws:

$$\begin{aligned}\bar{H}^1 &:= H^1 \left(\Omega^*(M^{(\infty)}) / \mathcal{I}^{(\infty)}, \bar{d} \right) \\ &= \frac{\{ \varphi \in \Omega^1(M^{(\infty)}) \mid d\varphi \in \mathcal{I}^{(\infty)} \}}{I^{(\infty)} + d\Omega^0(M^{(\infty)})}\end{aligned}$$

Analogy with Chern classes

$$\begin{array}{ccc}
 \mathbb{C}^k & \longrightarrow & V \\
 & & \downarrow \pi \\
 & & X
 \end{array}
 \quad \rightsquigarrow \quad
 c_i(V) \in H^{2i}(X, \mathbb{Z})$$

To each complex vector bundle we associate Chern classes that measure the *twisting of the bundle*.

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Universal Approach

The Chern classes of the universal bundle generate the cohomology of the infinite Grassmanian:

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Universal Approach

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The Chern classes of the universal bundle generate the cohomology of the infinite Grassmanian:

$$\begin{array}{ccc} \mathbb{C}^k & \longrightarrow & U_k \\ & & \downarrow \pi_k \\ & & Gr(k, \infty) \end{array}$$

Universal Approach

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathbb{Z} \cdot \langle c_1(U_k), \dots, c_k(U_k) \rangle = H^{2i}(Gr(k, \infty), \mathbb{Z})$$

Universal Approach

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathbb{Z} \cdot \langle c_1(U_k), \dots, c_k(U_k) \rangle = H^{2i}(Gr(k, \infty), \mathbb{Z})$$

$c_i(V) = f^*(c_i(U_k))$ for $f : X \rightarrow Gr(k, \infty)$ a classifying map.

Characteristic cohomology induces de Rham cohomology on solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions*Solution**EDS*

$$(N, \iota) \longrightarrow (M, \mathcal{I})$$

$$[\iota^*(\varphi)]_{dR} \in H_{dR}^1(N) \longleftarrow [\varphi]_{CC} \in \overline{H}^1$$

The space of generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
Cohomology**Structure
theorems for
 \overline{H}^1**

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

For $A : M^{(\infty)} \rightarrow \mathbb{C}$

The space of generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

For $A : M^{(\infty)} \rightarrow \mathbb{C}$

$$\mathcal{E}(A) := e_{-1} \overline{e_{-1}}(A) + f_u A.$$

The space of generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathcal{E}(A) := e_{-1} \overline{e_{-1}}(A) + f_u A.$$

$$V = \ker(\mathcal{E}) \cap \{P : P_{u_i \bar{u}_j} = P_u = 0\}$$

The space of generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$V_d = \left\{ P \in V \mid F_\lambda^*(P) = \lambda^d P \right\}$$

The space of generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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From now on, assume that $f_u \neq \alpha f + \beta$ for any $\alpha, \beta \in \mathbb{R}$.

The space of generating functions

Joint work with Oliver Goertsches (Selecta 2011)

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

The space of generating functions

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Theorem (A)

$$\textcircled{1} \quad V \cong \overline{H}^1$$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

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Theorem (A)

- ① $V \cong \overline{H}^1$
- ② $V = \bigoplus_{d \in \mathbb{Z}} V_d$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

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Theorem (A)

- ① $V \cong \overline{H}^1$
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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

Joint work with Oliver Goertsches (Selecta 2011)

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

Joint work with Oliver Goertsches (Selecta 2011)

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- ③ $\dim_{\mathbb{C}}(V_d) \leq 1$
- ④ $\dim_{\mathbb{C}}(V_{2n}) = 0$ for $n \neq 0$
- ⑤ For $d \neq 0$, $P \in V_d$ is a polynomial and $P = c \cdot u_{d-1} + \{u_0, u_1, \dots, u_{d-2}\}$, $c \in \mathbb{C}$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

Joint work with Oliver Goertsches (Selecta 2011)

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$$P = c \cdot u_{d-1} + \{u_0, u_1, \dots, u_{d-2}\}, \quad c \in \mathbb{C}$$
- ⑥ If f, f_u, f_{uu} are linearly independent over \mathbb{R} , then $V_d = 0$
 for $d \geq 2$: **No higher-order conservation laws**

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The space of generating functions

Joint work with Oliver Goertsches (Selecta 2011)

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$$P = c \cdot u_{d-1} + \{u_0, u_1, \dots, u_{d-2}\}, \quad c \in \mathbb{C}$$
- ⑥ If f, f_u, f_{uu} are linearly independent over \mathbb{R} , then $V_d = 0$
 for $d \geq 2$: **No higher-order conservation laws**
- ⑦ If $f_{uu} = \beta f$ for $\beta \in \mathbb{R}$ (e.g. $f = \sin$, $f = \sinh$) then
 $\dim_{\mathbb{C}}(V_{2n+1}) = 1$ for all $n \in \mathbb{Z}$.

Previous work

$$\frac{\partial^2 u}{\partial z \partial \bar{z}} = -f(u)$$

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- Olver (1977): Recursion operators involving $\frac{d}{dx}^{-1}$ to generate a hierarchy of conservation laws for $u_{xt} = \sin(u)$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Previous work

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Previous work

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Previous work

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Previous work

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Previous work

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Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Previous work

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

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Previous work

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

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Previous work

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Previous work

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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- These results do not directly imply a result about \overline{H}^1
- Is there a general formulation that will work for $u : \mathbb{C} \rightarrow \mathbb{R}^m$ —the Toda-field equations—when $m > 1$?

The Tzitzeica equation

Tzitzeica equation:

$$u_{z\bar{z}} = e^{-2u} - e^u \quad (f_{uu} = \alpha f_u + 2\alpha^2 f)$$

Goals

Nonlinear
Poisson
equation

Characteristic
Cohomology

Structure
theorems for
 \bar{H}^1

Previous work

**The Tzitzeica
equation**

Primitive
maps and
Killing fields

Killing fields
and
conservation
laws

The Killing
field recursion

Questions and
Conjectures

Finite type
solutions

The Tzitzeica equation

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$V_7 = \mathbb{C} \cdot \{u_6 + 7u_4u_1 - 7u_4u_0^2 + 14u_3u_2 - 28u_3u_1u_0 - 21u_2^2u_0 - 28u_2u_1^2 - 14u_2u_1u_0^2 + 14u_2u_0^4 - \frac{28}{3}u_1^3u_0 + 28u_1^2u_0^3 - \frac{4}{3}u_0^7\}$$

We want a recursion for generating functions of the characteristic cohomology when

$$f_{uu} = \alpha f_u + 2\alpha^2 f.$$

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Killing fields lead to a recursion that satisfies the first two and will very likely accommodate the third criterion

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

G compact real form of the semi-simple complex Lie group $G^{\mathbb{C}}$

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

G compact real form of the semi-simple complex Lie group $G^{\mathbb{C}}$
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Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$\tau : G^{\mathbb{C}} \rightarrow G^{\mathbb{C}}$ automorphism, $\tau^k = 1$, $\tau\sigma = \sigma\tau$, $K^{\mathbb{C}}$ the group it
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Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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stabilizes.

$(G^{\mathbb{C}}, \tau, \sigma)$ is a **k -symmetric space**

Eigen decomposition

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

σ, τ induce automorphisms of the Lie algebra $\mathfrak{g}^{\mathbb{C}}$ of $G^{\mathbb{C}}$.
(Assume that $\mathfrak{g}^{\mathbb{C}} \subset \mathfrak{gl}(r, \mathbb{C})$)

Eigen decomposition

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{g}_0 \oplus \mathfrak{g}_1 \oplus \dots \oplus \mathfrak{g}_{-1}$$

Eigen decomposition

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\mathfrak{g}^{\mathbb{C}} = \mathfrak{g}_0 \oplus \mathfrak{g}_1 \oplus \dots \oplus \mathfrak{g}_{-1}$$

where \mathfrak{g}_j is the eigenspace of τ with eigenvalue μ^j , for some
primitive k^{th} -root of unity μ .

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Let N be a simply connected Riemann surface.

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Let N be a simply connected Riemann surface. The map $\phi : N \rightarrow G/K$ has a framing $F : N \rightarrow G$.

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The Maurer-Cartan form ω on G induces a flat connection $\psi = F^*(\omega)$ on N .

Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Primitive maps

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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which is equivalent to

$$\psi_\lambda = \psi_{-1}\lambda^{-1} + \psi_0 + \psi_1\lambda \text{ is flat}$$

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Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Bolton, Pedit, Woodward (1995) show that, when $K = T^m$, there exist coordinates and a frame for which

$$F^{-1}F_z = \psi \left(\frac{\partial}{\partial z} \right) = U_z + Ad \exp(B)(U)$$

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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for $U : \mathbb{C} \rightarrow \mathbb{R}^m$ and $B \in \mathfrak{g}$ satisfying

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$U_{z\bar{z}} = -T(U) \quad (\text{Toda equations})$$

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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where $T(U) = \sum_{i,j} (A_i e^{a_{ij} U_j})$ is given in terms of roots of $\mathfrak{g}_{\mathbb{C}}$.

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

$SU(3)/SO(2)$ is a 6-symmetric space.

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$u_{z\bar{z}} = e^{-2u} - e^u$$

For the $SU(3)/SO(2)$ primitive map system, the Killing field equations can be turned into a recursion mapping $V_d \rightarrow V_{d+6}$.

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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There are two stages of the recursion that require a relative Poincare lemma.

Primitive maps and the Toda-field

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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There are two stages of the recursion that require a relative Poincare lemma. This relative Poincare lemma follows from the vanishing result $V_{2n} = 0$.

Killing fields

based loop algebra associated to $\mathfrak{g}^{\mathbb{C}}$

$$\mathcal{L}(\mathfrak{g}^{\mathbb{C}}) =$$

$$\left\{ B_{\lambda} \in \mathfrak{g}^{\mathbb{C}}[[\lambda, \lambda^{-1}]] : \text{if } B_{\lambda} = \sum_n B^n \lambda^n \text{ then } B^0 = 0 \right\}$$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Killing fields

based loop algebra associated to $\mathfrak{g}^{\mathbb{C}}$

$$\mathcal{L}(\mathfrak{g}^{\mathbb{C}}) =$$

$$\left\{ B_{\lambda} \in \mathfrak{g}^{\mathbb{C}}[[\lambda, \lambda^{-1}]] : \text{if } B_{\lambda} = \sum_n B^n \lambda^n \text{ then } B^0 = 0 \right\}$$

and the **twisted based loop algebra**

$$\mathcal{L}^{\sigma, \tau}(\mathfrak{g}^{\mathbb{C}}) = \left\{ B_{\lambda} \in \mathcal{L}(\mathfrak{g}^{\mathbb{C}}) : \sigma(B_{\lambda^{-1}}) = B_{\lambda} \quad \tau(B_{\mu\lambda}) = B_{\lambda} \right\}$$

associated to the k -symmetric space $(G^{\mathbb{C}}, \sigma, \tau)$.

Killing fields

based loop algebra associated to $\mathfrak{g}^{\mathbb{C}}$

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associated to the k -symmetric space $(G^{\mathbb{C}}, \sigma, \tau)$.

Definition

A **formal Killing field** for the family of flat connections ψ_{λ} on the Riemann surface N is a map $B_{\lambda} : N \rightarrow \mathcal{L}^{\sigma, \tau}(\mathfrak{g}^{\mathbb{C}})$ satisfying

$$dB_{\lambda} + [\psi_{\lambda}, B_{\lambda}] = 0.$$

Killing fields and infinitesimal symmetries

Burstall, Ferus, Pedit, Pinkall (1993) introduce Killing fields of Harmonic maps into symmetric spaces as a way to package infinitesimal symmetries (Jacobi fields)

Killing fields unpacked

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equation**Primitive
maps and
Killing fields**Killing fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The Killing field satisfies $B^{kn+j} \in \mathfrak{g}_j$.

Killing fields unpacked

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The Killing field satisfies $B^{kn+j} \in \mathfrak{g}_j$.

The Killing field equation decomposes into

$$\partial B^j + [\psi'_0, B^j] + [\psi_{-1}, B^{j+1}] = 0$$

$$\bar{\partial} B^j + [\psi''_0, B^j] + [\psi_1, B^{j-1}] = 0.$$

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fields**Killing fields
and
conservation
laws**The Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Assume that K is abelian.

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fields**Killing fields
and
conservation
laws**The Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Assume that K is abelian.

Let $A_{-1} = \psi_{-1} \left(\frac{\partial}{\partial z} \right)$, $A_1 = \psi_1 \left(\frac{\partial}{\partial \overline{z}} \right)$ and

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fields**Killing fields
and
conservation
laws**The Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Assume that K is abelian.

Let $A_{-1} = \psi_{-1} \left(\frac{\partial}{\partial z} \right)$, $A_1 = \psi_1 \left(\frac{\partial}{\partial \bar{z}} \right)$ and

For $P : N \rightarrow \mathfrak{g}_0$ define

$$\mathcal{E}_{(G^{\mathbb{C}}, \sigma, \tau)}(P) = \Delta P + 4[A_{-1}, [A_1, P]].$$

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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For $P : N \rightarrow \mathfrak{g}_0$ define

$$\mathcal{E}_{(G^{\mathbb{C}}, \sigma, \tau)}(P) = \Delta P + 4[A_{-1}, [A_1, P]].$$

We find that $\mathcal{E}_{(G^{\mathbb{C}}, \sigma, \tau)}(B^{kn}) = 0$ for $B^{kn} \in \mathfrak{g}_0$.

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fields**Killing fields
and
conservation
laws**The Killing
field recursionQuestions and
ConjecturesFinite type
solutions

For any pair of functions $P, Q \in \Omega^0(N, \mathfrak{g}_0)$ define

$$\varphi_{P,Q} = -\sqrt{-1}J(\kappa(P, dQ) - \kappa(Q, dP)) \in \Omega^1(N, \mathbb{C})$$

where κ is the Killing form of $\mathfrak{g}^{\mathbb{C}}$.

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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where κ is the Killing form of $\mathfrak{g}^{\mathbb{C}}$.

Lemma

The one-form $\varphi_{P,Q}$ is closed if $P, Q \in \ker(\mathcal{E}_{(G^{\mathbb{C}}, \sigma, \tau)})$.

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The \mathfrak{g}_0 -components of Killing fields give rise to conservation laws for primitive maps (when $K = T^m$).

Killing fields and conservation laws

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The one-form $\varphi_{P,Q}$ is closed if $P, Q \in \ker(\mathcal{E}_{(G^{\mathbb{C}}, \sigma, \tau)})$.

The \mathfrak{g}_0 -components of Killing fields give rise to conservation laws for primitive maps (when $K = T^m$).

Terng and Wang (2004) gave a similar formula for conservation laws of the U/K -systems using Killing field-like objects.

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

The Recursion: Let $a^n \in V_{2n+1}$ for some $n \geq 0$.

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

The Recursion: Let $a^n \in V_{2n+1}$ for some $n \geq 0$.

Define the one-form

$$\alpha^n = \frac{\sqrt{-1}}{3\sqrt{2}}(a_{-1,-1}^n + 2u_0 a_{-1}^n)\zeta - \frac{\sqrt{-1}}{\sqrt{2}}e^u a^n \bar{\zeta} \quad (1)$$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$d\alpha^n \equiv 0 \text{ modulo } I^{(\infty)}$$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$d\alpha^n \equiv 0$ modulo $I^{(\infty)}$ so $[\alpha^n] \in \overline{H}^1$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$d\alpha^n \equiv 0$ modulo $I^{(\infty)}$ so $[\alpha^n] \in \overline{H}^1$

Thus it must correspond to an element of $P_\alpha \in V_{2n+2} = 0$.

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Thus $[\alpha^n] = 0$

There exists $b^n : M^{(\infty)} \rightarrow \mathbb{C}$ of weighted-degree $2n + 2$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Thus $[\alpha^n] = 0$

There exists $b^n : M^{(\infty)} \rightarrow \mathbb{C}$ of weighted-degree $2n+2$ such that $db^n \equiv \alpha^n$ modulo $I^{(\infty)}$.

The recursion for $f_{uu} = \alpha f_u + \beta f$

The Recursion: Let $a^n \in V_{2n+1}$ for some $n \geq 0$.

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$$\alpha^n = \frac{\sqrt{-1}}{3\sqrt{2}}(a_{-1,-1}^n + 2u_0 a_{-1}^n)\zeta - \frac{\sqrt{-1}}{\sqrt{2}}e^u a^n \bar{\zeta} \quad (1)$$

$d\alpha^n \equiv 0$ modulo $I^{(\infty)}$ so $[\alpha^n] \in \overline{H}^1$

Thus it must correspond to an element of $P_\alpha \in V_{2n+2} = 0$.

Thus $[\alpha^n] = 0$

There exists $b^n : M^{(\infty)} \rightarrow \mathbb{C}$ of weighted-degree $2n+2$ such that $db^n \equiv \alpha^n$ modulo $I^{(\infty)}$.

b^n is a polynomial in u_0, \dots, u_{2n+1} .

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Recursively construct

$$f^n = \sqrt{-1} e^{\frac{u}{2}} (b_{-1}^n - u_0 b^n)$$

$$r^n = \frac{1}{3\sqrt{2}} e^{-\frac{u}{2}} (f_{-1}^n + \frac{1}{2} u_0 f^n)$$

$$s^n = -\frac{1}{\sqrt{2}} e^{-\frac{u}{2}} r_{-1}^n.$$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Again, define a one-form

$$\beta^n = \frac{\sqrt{-1}}{3} e^u (s_{-1,-1}^n - u_0 s_{-1}^n) \zeta - \sqrt{-1} e^{-u} s^n \bar{\zeta} \quad (2)$$

Once again, $d\beta^n \equiv 0$ modulo $I^{(\infty)}$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Again, define a one-form

$$\beta^n = \frac{\sqrt{-1}}{3} e^u (s_{-1,-1}^n - u_0 s_{-1}^n) \zeta - \sqrt{-1} e^{-u} s^n \bar{\zeta} \quad (2)$$

Once again, $d\beta^n \equiv 0$ modulo $I^{(\infty)}$

$$V_{2n+6} = 0$$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Again, define a one-form

$$\beta^n = \frac{\sqrt{-1}}{3} e^u (s_{-1,-1}^n - u_0 s_{-1}^n) \zeta - \sqrt{-1} e^{-u} s^n \bar{\zeta} \quad (2)$$

Once again, $d\beta^n \equiv 0$ modulo $I^{(\infty)}$

$V_{2n+6} = 0$ implies there exists $t^n : M^{(\infty)} \rightarrow \mathbb{C}$ of weighted-degree $2n + 6$ such that $dt^n \equiv \beta^n$ modulo $I^{(\infty)}$.

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Again, define a one-form

$$\beta^n = \frac{\sqrt{-1}}{3} e^u (s_{-1,-1}^n - u_0 s_{-1}^n) \zeta - \sqrt{-1} e^{-u} s^n \bar{\zeta} \quad (2)$$

Once again, $d\beta^n \equiv 0$ modulo $I^{(\infty)}$

$V_{2n+6} = 0$ implies there exists $t^n : M^{(\infty)} \rightarrow \mathbb{C}$ of weighted-degree $2n+6$ such that $dt^n \equiv \beta^n$ modulo $I^{(\infty)}$.

t^n is a polynomial in u_0, \dots, u_{2n+5} .

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Finally define

$$a^{n+1} = -\sqrt{-2}(t_{-1}^n + u_0 t^n) \quad (3)$$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Finally define

$$a^{n+1} = -\sqrt{-2}(t_{-1}^n + u_0 t^n) \quad (3)$$

One checks that $a^{n+1} \in V_{2n+7}$

The recursion for $f_{uu} = \alpha f_u + \beta f$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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One checks that $a^{n+1} \in V_{2n+7}$

Summary: There is a level six recursion derived from the Killing field equation; requires $V_{2n} = 0$.

Generating functions

Using the seeds

$$V_1 = \mathbb{C} \cdot \{u_0\}$$

Generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

Using the seeds

$$V_1 = \mathbb{C} \cdot \{u_0\}$$

$$V_5 = \mathbb{C} \cdot \{u_4 + 5u_2u_1 - 5u_2u_0^2 - 5u_1^2u_0 + u_0^5\}$$

Generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
laws**The Killing
field recursion**Questions and
ConjecturesFinite type
solutions

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$$V_9 = \mathbb{C} \cdot \{u_8 + \cdots\}$$

Generating functions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Using the seeds

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$$V_5 = \mathbb{C} \cdot \{u_4 + 5u_2u_1 - 5u_2u_0^2 - 5u_1^2u_0 + u_0^5\}$$

$$V_9 = \mathbb{C} \cdot \{u_8 + \cdots\}$$

for the order six recursion, implies that $V_{2n+1} \cong \mathbb{C}$ for $n \in \mathbb{Z}$
and $n \neq 1, -2$.

Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Determine the characteristic cohomology for all Toda-field equations associated with primitive maps to k -symmetric spaces.

Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Determine the characteristic cohomology for all Toda-field equations associated with primitive maps to k -symmetric spaces.

We have done the cases in which $K = \mathrm{SO}(2)$:

Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Determine the characteristic cohomology for all Toda-field equations associated with primitive maps to k -symmetric spaces.

We have done the cases in which $K = \mathrm{SO}(2)$:

- 1 $\mathrm{SU}(2)/\mathrm{SO}(2)$ —Gauss maps of CMC surfaces in \mathbb{R}^3

Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Determine the characteristic cohomology for all Toda-field equations associated with primitive maps to k -symmetric spaces.

We have done the cases in which $K = \mathrm{SO}(2)$:

- 1 $\mathrm{SU}(2)/\mathrm{SO}(2)$ —Gauss maps of CMC surfaces in \mathbb{R}^3
2-symmetric space

Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Determine the characteristic cohomology for all Toda-field equations associated with primitive maps to k -symmetric spaces.

We have done the cases in which $K = \mathrm{SO}(2)$:

- ① $\mathrm{SU}(2)/\mathrm{SO}(2)$ —Gauss maps of CMC surfaces in \mathbb{R}^3 2-symmetric space
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Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Toda-field equations

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Finite type solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Let $\iota : N \rightarrow (M, \mathcal{I})$ be an integral manifold of the EDS associated to the nonlinear Poisson equation.

Finite type solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Let $\bar{H}^1 = \mathbb{R} \cdot \{[\varphi_1], [\varphi_2], \dots\}$

Finite type solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Let $\bar{H}^1 = \mathbb{R} \cdot \{[\varphi_1], [\varphi_2], \dots\}$ and let $\mathcal{H} = \mathbb{R} \cdot \{\varphi_1, \varphi_2, \dots\}$ with φ_i in ‘normal form.’

Finite type solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Definition

The integral manifold $\iota : N \rightarrow M$ is of finite type if $\dim_{\mathbb{R}}(\iota^*(\mathcal{H})) < \infty$.

Finite type solutions

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The integral manifold $\iota : N \rightarrow M$ is of finite type if $\dim_{\mathbb{R}}(\iota^*(\mathcal{H})) < \infty$.

This agrees with the definition of finite type given by Pinkall and Sterling.

A linear embedding—the spectral curve?

Suppose that $\iota : \mathbb{C} \rightarrow M^{(\infty)}$ is a doubly periodic solution for the lattice $\Lambda \subset \mathbb{C}$, of finite type n .

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

A linear embedding—the spectral curve?

Suppose that $\iota : \mathbb{C} \rightarrow M^{(\infty)}$ is a doubly periodic solution for the lattice $\Lambda \subset \mathbb{C}$, of finite type n .

Let $\{[\varphi_j]\}$ be a basis for \overline{H}^1

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Let $\{[\varphi_j]\}$ be a basis for \overline{H}^1 such that the generating function P_j of φ_j satisfies $P_j = u_{j-1} + \cdots$ for $j \geq 1$.

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \overline{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Let $\bar{H}_n^1 := \{[\varphi_j] \in \bar{H}^1 \mid j \leq n\} \cong \mathbb{R}^{2n}$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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$$\begin{aligned} \iota^* : \bar{H}_n^1 &\rightarrow H^1(\mathbb{C}/\Lambda, \mathbb{R}) \\ \mathbb{R}^{2n} &\rightarrow \mathbb{R}^2 \end{aligned}$$

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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A linear embedding—the spectral curve?

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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The angles determining this linear embedding should contain geometric information.

A linear embedding—the spectral curve?

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

Conjecture

$\hat{\iota}^* : H_1(\mathbb{C}/\Lambda, \mathbb{R}) \rightarrow (\bar{H}_n^1)^*$ *descends to a map*

$$\hat{\iota}^* : H_1(\mathbb{C}/\Lambda, \mathbb{R})/\Lambda^* \rightarrow \mathbb{R}^{2n}/\Gamma \subset \text{Jac}(X_u)$$

where $\Gamma \subset \mathbb{R}^{2n}$ is a lattice of real rank $2n$,

A linear embedding—the spectral curve?

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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A linear embedding—the spectral curve?

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Idea: Conservation laws define an *extended* Abel-Jacobi map

Global approach

Global conservation laws for minimal surfaces $N \rightarrow \mathbb{S}^3$ might lead to a maximally linear embedding

$$N \rightarrow \text{Jac}(N) \hookrightarrow T^{2n}$$

even if N is a compact Riemann surface with genus > 1 , thus extending the spectral curve type construction to higher genus domains.

Goals

Nonlinear
Poisson
equationCharacteristic
CohomologyStructure
theorems for
 \bar{H}^1

Previous work

The Tzitzeica
equationPrimitive
maps and
Killing fieldsKilling fields
and
conservation
lawsThe Killing
field recursionQuestions and
ConjecturesFinite type
solutions

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Likely to require:

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Pedit has suggested the same possibility based on the multiplier spectral curve.